

FRIEDRICHS' EXTENSION LEMMA WITH BOUNDARY VALUES AND APPLICATIONS IN COMPLEX ANALYSIS

J. RUPPENTHAL

ABSTRACT. Let Q be a first-order differential operator on a compact, smooth oriented Riemannian manifold with smooth boundary. Then, Friedrichs' extension lemma states that the minimal closed extension Q_{min} (the closure of the graph) and the maximal closed extension Q_{max} (in the sense of distributions) of Q in L^p -spaces ($1 \leq p < \infty$) coincide. In the present paper, we show that the same is true for boundary values with respect to Q_{min} and Q_{max} . This gives a useful characterization of weak boundary values, particularly for $Q = \bar{\partial}$ the Cauchy-Riemann operator. As an application, we derive the Bochner-Martinelli-Koppelman formula for L^p -forms with weak $\bar{\partial}$ -boundary values.

1. INTRODUCTION

Let D be a relatively compact domain in a Hermitian complex manifold and $\bar{\partial} : C_*^\infty(\bar{D}) \rightarrow C_*^\infty(\bar{D})$ the Cauchy-Riemann operator on smooth forms. For $1 \leq p < \infty$, this operator can be considered as a densely defined graph-closable operator on L^p -forms:

$$\bar{\partial} : \text{Dom}(\bar{\partial}) = C_*^\infty(\bar{D}) \subset L_*^p(D) \rightarrow L_*^p(D)$$

Now then, the $\bar{\partial}$ -operator has various closed extensions. The two most important are the minimal closed extension $\bar{\partial}_{min}$ given by the closure of the graph and the maximal closed extension $\bar{\partial}_{max}$, i.e. the $\bar{\partial}$ -operator in the sense of distributions. Whereas the two extensions coincide on smoothly bounded domains by Friedrichs' extension lemma (see [F], [H3]), one has to be very careful when considering non-smooth domains. Especially on regular sets in singular complex spaces, it is crucial to distinguish the different closed extensions of the $\bar{\partial}$ -operator for they lead to different Dolbeault cohomology groups (see e.g. [BS], [P], [PS1] or [PS2]). It was realized that investigating the relation between the various extensions is an essential and very fruitful (maybe even indispensable) step in understanding the $\bar{\partial}$ -equation on singular complex spaces which has to be pursued (see also [R2]). Clearly, the difference between the closed extensions occurs at the boundary of the domain. So, a first step is to study the boundary behavior of $\bar{\partial}_{min}$ and $\bar{\partial}_{max}$ on domains with smooth boundary which we do in the present paper by deriving a Friedrichs' extension lemma with boundary values.

Date: October 14, 2009.

2000 Mathematics Subject Classification. 32W05, 35F99, 47F05.

Key words and phrases. Weak boundary values, generalized Stokes' formulas, Cauchy-Riemann operator.

Let $D \subset\subset \mathbb{C}^n$ be a bounded domain with smooth boundary bD , and let $f \in L^p_{0,q}(D)$ with $\bar{\partial}f \in L^p_{0,q+1}(D)$ in the sense of distributions for $1 \leq p < \infty$. Then, we say that f has weak $\bar{\partial}$ -boundary values $f_b \in L^p_q(bD)$ in the sense of distributions if

$$\int_D \bar{\partial}f \wedge \phi + (-1)^q \int f \wedge \bar{\partial}\phi = \int_{bD} f_b \wedge \iota^*(\phi) \quad (1)$$

for all $\phi \in C^\infty_{n,n-q-1}(\bar{D})$, where $\iota : bD \hookrightarrow \mathbb{C}^n$ is the embedding of the boundary. Weak $\bar{\partial}$ -boundary values in the sense of distributions are a classical subject of complex analysis (see Theorem 4.2, for example) and closely related to the investigation of the so-called Hardy spaces (cf. [S]). Starting from results of Skoda [S], Harvey and Polking [HP], Schuldenzucker [SCH] and Hefer [H1], there has been a considerable progress in the understanding of weak $\bar{\partial}$ -boundary values by Hefer in [H2], where boundary values in the sense of distributions are compared to boundary values which arise naturally in the application of integral operators. This is interesting because boundary values defined by restricting the kernel of an integral operator can often be estimated by direct methods, whereas the abstractly given distributional boundary values are less tractable but analytically interesting objects linked to the form on the interior of a domain.

However, in applications the definition of weak $\bar{\partial}$ -boundary values by means of the Stokes' formula (1) turns out to be a bit unhandy and it is more convenient to have boundary values in the sense of approximation by smooth forms. In fact, let $f \in \text{Dom}(\bar{\partial}_{max}) \subset L^p_{0,q}(D)$ with weak boundary values $f_b \in L^p_q(bD)$ according to definition (1), and let $r \in C^\infty(\mathbb{C}^n)$ be a smooth defining function for D . Then we will show that there exists a sequence $f_j \in C^\infty_{0,q}(\bar{D})$ such that

$$f_j \rightarrow f \text{ in } L^p_{0,q}(D), \quad \bar{\partial}f_j \rightarrow \bar{\partial}f \text{ in } L^p_{0,q+1}(D) \quad (2)$$

(the classical Friedrichs' extension lemma) and moreover

$$f_j \wedge \partial r \rightarrow f_b \wedge \partial r \text{ on } bD \text{ in } L^p_q(bD), \quad (3)$$

i.e. f has $\bar{\partial}$ -boundary values in the sense of approximation (Theorem 4.4).

This phenomenon is not restricted to the Cauchy-Riemann operator, but holds for arbitrary differential operators of first order with smooth coefficients. So, it is more convenient to adopt a more general point of view. Let M be a smooth, compact Riemannian manifold with smooth boundary, E and F Hermitian vector bundles over M , and $Q : C^\infty(M, E) \rightarrow C^\infty(M, F)$ a differential operator of first order. Let $1 \leq p < \infty$ and $f \in L^p(M, E)$. We say that $f \in \text{Dom}(Q^p_{min})$ if there exists a sequence $\{f_j\} \subset C^\infty(M, E)$ and a section $g \in L^p(M, F)$ such that

$$f_j \rightarrow f \text{ in } L^p(M, E), \quad Qf_j \rightarrow g \text{ in } L^p(M, F),$$

and define $Q^p_{min}f := g$. The well-defined operator Q^p_{min} is called the minimal extension of Q because it is the closed extension of Q to an operator $L^p(M, E) \rightarrow L^p(M, F)$ with minimal domain of definition. Its graph is simply

the closure of the graph of $Q : C^\infty(M, E) \rightarrow C^\infty(M, F)$ in $L^p(M, E) \times L^p(M, F)$. Let σ_Q be the principal symbol of Q , ν the outward pointing unit normal to bM , and $\nu^\flat$ the dual cotangent vector. Then, we say that f has boundary values with respect to Q_{min}^p if there exists a sequence $\{f_j\}$ in $C^\infty(M, E)$ such that $\lim_{j \rightarrow \infty} f_j = f$ in $L^p(M, E)$, $\lim_{j \rightarrow \infty} Qf_j = Q_{min}^p f$ in $L^p(M, F)$, and a section $f_b \in L^p(bM, E|_{bM})$ such that

$$\lim_{j \rightarrow \infty} \sigma_Q(\cdot, \nu^\flat(\cdot))f_j|_{bM} = \sigma_Q(\cdot, \nu^\flat(\cdot))f_b \quad \text{in } L^p(bM, F|_{bM}).$$

In this case, we call f_b weak Q -boundary values of f with respect to Q_{min}^p (i.e. in the sense of approximation).

Now, we draw our attention to the maximal closed extension of Q , that is the extension of Q in the sense of distributions. We say that $f \in \text{Dom}(Q_{max}^p)$ if $Qf = u \in L^p(M, F)$ in the sense of distributions, and set $Q_{max}^p f := u$ in that case. Here again, we can define weak Q -boundary values with respect to Q_{max}^p . We say that f has weak Q -boundary values $f_b \in L^p(bM, E|_{bM})$ with respect to Q_{max}^p (in the sense of distributions), if f_b satisfies the generalized Green-Stokes formula (cf. Theorem 2.1)

$$(Qu, \phi)_M - (u, Q^* \phi)_M = \frac{1}{i} \int_{bM} \langle \sigma_Q(x, \nu^\flat) u_b, \phi \rangle_{F_x} dS(x)$$

for all $\phi \in C^\infty(M, F)$.

The main objective of the present paper is to compare both notions of Q -boundary values. It is easy to see that $\text{Dom}(Q_{min}^p) \subset \text{Dom}(Q_{max}^p) \subset L^p(M, E)$, and that Q_{min}^p is the restriction of Q_{max}^p to $\text{Dom}(Q_{min}^p)$. Moreover, it is also clear that weak Q -boundary values in the sense of approximation are weak Q -boundary values in the sense of distributions as well. It is well-known that in fact $Q_{min}^p = Q_{max}^p$ on smooth, compact manifolds with smooth boundary. This result, due to Friedrichs (see [F], [H3]), is usually called Friedrichs' extension lemma (Theorem 3.1). In this paper, we observe that the two notions of boundary values coincide as well (Theorem 3.3). One might call this Friedrichs' extension lemma with boundary values. In the particular case of the Cauchy-Riemann operator $Q = \bar{\partial}$, we obtain (2), (3).

The present paper is organized as follows: in the next section, we recall the notion of weak Q -boundary values in the sense of distributions (Definition 2.2) which makes sense in view of the generalized Green-Stokes formula Theorem 2.1. In section 3, we recall the proof of the classical Friedrichs' extension lemma as it is presented in [LM] (relying on [H3] which in turn cites [F] and [LP]) and prove Friedrichs' extension lemma with boundary values by a sophisticated choice of an approximating identity. In section 4, we return to the Cauchy-Riemann operator by translating the results into the language of complex analysis in the particular case of the differential operator $Q = \bar{\partial}$. In the last section, we show how boundary values in the sense of approximation can be used in applications by deriving the Bochner-Martinelli-Koppelman formula for forms with weak $\bar{\partial}$ -boundary values.

2. WEAK BOUNDARY VALUES

Let M be a smooth, compact Riemannian manifold with smooth boundary, E and F Hermitian vector bundles over M , and

$$Q : C^\infty(M, E) \rightarrow C^\infty(M, F)$$

a differential operator of first order. Let σ_Q be the principal symbol of Q , and

$$Q^* : C^\infty(M, F) \rightarrow C^\infty(M, E)$$

its formal adjoint operator given by

$$(Qu, v)_M = \int_M \langle Qu, v \rangle_F dV_M = \int_M \langle u, Q^*v \rangle_E dV_M = (u, Q^*v)_M,$$

where one of the two sections $u \in C^\infty(M, E)$, $v \in C^\infty(M, F)$ has compact support in the interior of M . Let dS be the induced volume element on the boundary bM , ν the outward pointing unit normal to bM , and $\nu^\flat$ the dual cotangent vector. Then, the generalized Green-Stokes formula reads as (see [T], Prop. 9.1):

Theorem 2.1. *Let M be a smooth, compact Riemannian manifold with smooth boundary, and Q a first-order differential operator (acting on sections of Hermitian vector bundles). Then*

$$(Qu, v)_M - (u, Q^*v)_M = \frac{1}{i} \int_{bM} \langle \sigma_Q(x, \nu^\flat)u, v \rangle_{F_x} dS(x)$$

for all sections $u \in C^\infty(M, E)$, $v \in C^\infty(M, F)$.

Now, let $u \in L^1(M, E)$ and $f \in L^1(M, F)$. Then we say that $Qu = f$ in the sense of distributions if

$$(u, Q^*\phi)_M = (f, \phi)_M$$

for all $\phi \in C^\infty(M, F)$ with compact support in the interior of M . We can now give the definition of weak boundary values with respect to the first-order differential operator Q :

Definition 2.2. *In the situation of Theorem 2.1, let $u \in L^1(M, E)$ with $Qu \in L^1(M, F)$. Then u has weak Q -boundary values $u_b \in L^p(bM, E|_{bM})$ if*

$$(Qu, \phi)_M - (u, Q^*\phi)_M = \frac{1}{i} \int_{bM} \langle \sigma_Q(x, \nu^\flat)u_b, v \rangle_{F_x} dS(x) \quad (4)$$

for all $\phi \in C^\infty(M, F)$.

This generalizes the notion of weak boundary values of functions in the Sobolev space $H^{1,p}(M)$: Let $Q = d : C^\infty(M, \mathbb{C}) \rightarrow C^\infty(M, \mathbb{C} \otimes T^*M)$ be the exterior derivative. Then, for all $1 \leq p \leq \infty$, there is a unique continuous trace operator

$$T : H^{1,p}(M) = \{u \in L^p(M, \mathbb{C}) : du \in L^p(M, \mathbb{C} \otimes T^*M)\} \rightarrow L^p(bM, \mathbb{C})$$

such that Tu satisfies (4) (cf. [A], A 6.6). In general, weak Q -boundary values do not necessarily exist.

3. FRIEDRICHS' EXTENSION LEMMA

Again, let M be a smooth, compact Riemannian manifold with smooth boundary, E and F Hermitian vector bundles over M , and $Q : C^\infty(M, E) \rightarrow C^\infty(M, F)$ a differential operator of first order (with C^1 coefficients). Let $1 \leq p < \infty$. Then, for $f \in L^p(M, E)$, we say that $f \in \text{Dom}(Q_{min}^p)$ if there exists a sequence $\{f_j\} \subset C^\infty(M, E)$ and $g \in L^p(M, F)$ such that

$$f_j \rightarrow f \text{ in } L^p(M, E), \quad Qf_j \rightarrow g \text{ in } L^p(M, F),$$

and define $Q_{min}^p f := g$. The operator Q_{min}^p is uniquely defined, because

$$(g, h)_M = \lim_{j \rightarrow \infty} (Qf_j, h)_M = \lim_{j \rightarrow \infty} (f_j, Q^*h)_M = (f, Q^*h)_M$$

for all $h \in C^\infty(M, F)$ with compact support in the interior of M . Moreover, we say that $f \in \text{Dom}(Q_{max}^p)$, if $Qf = u \in L^p(M, F)$ in the sense of distributions, and set $Q_{max}^p f := u$ in that case. It is easy to see that

$$\text{Dom}(Q_{min}^p) \subset \text{Dom}(Q_{max}^p) \subset L^p(M, E),$$

and Q_{min}^p is the restriction of Q_{max}^p to $\text{Dom}(Q_{min}^p)$. But, in our situation, also the converse is true (cf. [LM], Theorem V.2.6):

Theorem 3.1. (Friedrichs' Extension Lemma) *Let M be a smooth, compact Riemannian manifold with smooth boundary, and $Q : C^\infty(M, E) \rightarrow C^\infty(M, F)$ a first-order differential operator (acting on sections of Hermitian vector bundles), and $1 \leq p < \infty$. Then for any $f \in \text{Dom}(Q_{max}^p)$ there exists a sequence $\{f_\epsilon\}$ in $C^\infty(M, E)$ such that $\lim_{\epsilon \rightarrow 0} f_\epsilon = f$ and $\lim_{\epsilon \rightarrow 0} Qf_\epsilon = Q_{max}^p f$ with respect to L^p -norms. Shortly this means that*

$$Q_{min}^p = Q_{max}^p.$$

Let us recall the principles of the proof. Using a partition of unity, it is enough to consider $U \subset\subset \mathbb{R}^n$ open with smooth boundary and $Q : C^\infty(U) \rightarrow C^\infty(U)$. So, let $f \in L^p(U)$ and $Qf = Q_{max}^p f \in L^p(U)$. Again, by the partition of unity argument, one has to consider the following two cases:

1. $\text{supp}(f) \subset\subset U$, or
2. $U = \{x \in \mathbb{R}^n : x_1 < 0\}$ and $\text{supp}(f) \subset\subset \overline{U}$.

For the first case, let $\phi \in C_{cpt}^\infty(B_1(0))$ with $\phi \geq 0$ and $\int \phi dx = 1$, where dx is the Euclidean volume element. We call $\phi_\epsilon(x) := \epsilon^{-n} \phi(x/\epsilon)$ a Dirac sequence, and

$$f_\epsilon := f * \phi_\epsilon$$

the convolution of f with a Dirac sequence. It is well known that $f_\epsilon \rightarrow f$ in $L^p(U)$ for $\epsilon \rightarrow 0^+$. But the crucial observation is

Lemma 3.2.

$$\|Qf_\epsilon - (Qf) * \phi_\epsilon\|_{L^p(U)} \lesssim \|f\|_{L^p(U)}.$$

It is now easy to complete the first case: Let $\delta > 0$ and $\psi \in C_{cpt}^\infty(U)$ such that

$$\|f - \psi\|_{L^p(U)} < \delta.$$

Applying Lemma 3.2 to $f - \psi$ yields:

$$\|Qf_\epsilon - (Qf) * \phi_\epsilon\|_{L^p(U)} \lesssim \delta + \|Q\psi_\epsilon - (Q\psi) * \phi_\epsilon\|_{L^p(U)}.$$

Choosing δ and ϵ arbitrarily small finishes this part of the proof. The second case is treated by exactly the same procedure. One only has to be a little careful when choosing the Dirac sequence ϕ_ϵ . Here, let $\phi \in C_{cpt}^\infty(B_1(0))$ such that

$$\text{supp}(\phi) \subset\subset \{x \in B_1(0) : x_1 > 0\}.$$

Then f_ϵ is well defined on $\bar{U}$, Lemma 3.2 is still true and everything goes through as before. That completes the proof of Theorem 3.1 as it is given in [LM].

We are now interested in the behavior of the sequence $\{f_\epsilon\}$ on the boundary bM . It is possible to extend Theorem 3.1 to Friedrichs' extension lemma with boundary values:

Theorem 3.3. *In the situation of Theorem 3.1, assume that $f \in \text{Dom}(Q_{max}^p)$ has weak Q -boundary values $f_b \in L^p(bM, E|_{bM})$ in the sense of Definition 2.2. Then there exists a sequence $\{f_\epsilon\}$ in $C^\infty(M, E)$ such that $\lim_{\epsilon \rightarrow 0} f_\epsilon = f$ in $L^p(M, E)$, $\lim_{\epsilon \rightarrow 0} Qf_\epsilon = Q_{max}^p f$ in $L^p(M, F)$ and*

$$\lim_{\epsilon \rightarrow 0} \sigma_Q(\cdot, \nu^\flat(\cdot))f_\epsilon|_{bM} = \sigma_Q(\cdot, \nu^\flat(\cdot))f_b \quad \text{in } L^p(bM, F|_{bM}).$$

Proof. We copy the proof of Theorem 3.1. One has to be even more careful when choosing the Dirac sequence. We only have to take a closer look at the second case. So, let $U = \{x \in \mathbb{R}^n : x_1 < 0\}$, $\text{supp}(f) \subset\subset \bar{U}$ and

$$\text{supp}(f_b) \subset\subset bU = \{x \in \mathbb{R}^n : x_1 = 0\}.$$

Then

$$Q = \sum_{j=1}^n a_j(x) \frac{\partial}{\partial x_j} + b(x), \quad Q^* = - \sum_{j=1}^n \left(a_j(x) \frac{\partial}{\partial x_j} + \frac{\partial a_j}{\partial x_j}(x) \right) + b(x),$$

and

$$\int_U (Qf)\Phi dx - \int_U f(Q^*\Phi) dx = \int_{bU} a_1(0, x') f_b(x') \Phi(0, x') dx' \quad (5)$$

for all $\Phi \in C_{cpt}^\infty(\bar{U})$ according to Definition 2.2 with $x' = (x_2, \dots, x_n)$. We will use the decomposition

$$Q^* = -a_1(x) \frac{\partial}{\partial x_1} + Q'. \quad (6)$$

Now, let us choose the right Dirac sequence for our purposes. Let $B'_1(0)$ be the unit ball in $\mathbb{R}^{n-1}$ and $\psi \in C_{cpt}^\infty(B'_1(0))$ with $\psi \geq 0$ and $\int \psi dx' = 1$, where dx' is the Euclidean volume element in $\mathbb{R}^{n-1}$, $x' = (x_2, \dots, x_n)$. For $\epsilon > 0$ set

$$\psi_\epsilon := \epsilon^{-(n-1)} \psi(x'/\epsilon).$$

Therefore, it follows that

$$\lim_{\epsilon \rightarrow 0} (a_1(0, \cdot) f_b) * \psi_\epsilon = a_1(0, \cdot) f_b \quad \text{in } L^p(bU). \quad (7)$$

Moreover, let $h : \mathbb{R} \rightarrow [0, 1]$ be a smooth function such that

$$h(x_1) = \begin{cases} 0, & \text{for } x_1 \leq 1, \\ 1, & \text{for } x_1 \geq 2. \end{cases}$$

For $\tau > 0$, set $h_\tau(x_1) = h(x_1/\tau)$. Now, we define a Dirac sequence in $\mathbb{R}^n$:

$$\phi_\epsilon(x) := \psi_\epsilon(x') \frac{\partial h_{\tau(\epsilon)}}{\partial x_1}(x_1),$$

where $\tau(\epsilon)$ will be chosen later. At this point, we only require that $\tau(\epsilon) \leq \epsilon$. Note that $\text{supp}(\phi_\epsilon) \subset \subset \{x_1 > 0\}$. Let

$$f_\epsilon := f * \phi_\epsilon.$$

Then $f_\epsilon \rightarrow f$ and $Qf_\epsilon \rightarrow Qf$ as in the proof of Theorem 3.1. Because of (7), we only have to prove that

$$\lim_{\epsilon \rightarrow 0} a_1(0, \cdot) f_\epsilon|_{bU} = \lim_{\epsilon \rightarrow 0} (a_1(0, \cdot) f_b) * \psi_\epsilon \quad \text{in } L^p(bU). \quad (8)$$

For $(0, x') \in bU$, we calculate:

$$\begin{aligned} (a_1(0, \cdot) f_b) * \psi_\epsilon(x') &= \int_{bU} a_1(0, t') f_b(t') \psi_\epsilon(x' - t') dt' \\ &= \int_{bU} a_1(0, t') f_b(t') \psi_\epsilon(x' - t') (1 - h_{\tau(\epsilon)}(0)) dt' \\ &= \int_U (Qf) \Phi_{x'}^\epsilon dt - \int_U f(Q^* \Phi_{x'}^\epsilon) dt, \end{aligned}$$

if we let

$$\Phi_{x'}^\epsilon(t) := \psi_\epsilon(x' - t') (1 - h_{\tau(\epsilon)}(-t_1))$$

and apply the Green-Stokes formula (5). By the use of the decomposition (6), it follows that

$$\begin{aligned} (a_1(0, \cdot) f_b) * \psi_\epsilon(x') &= \int_U (Qf) \Phi_{x'}^\epsilon dt - \int_U f(Q' \Phi_{x'}^\epsilon) dt \\ &\quad + \int_U f(t) a_1(t) \phi_\epsilon(x - t) dt \end{aligned}$$

with $x = (0, x')$. We will now show that the first two terms turn to 0 in $L^p(bU)$ if we let $\epsilon \rightarrow 0$.

So, consider:

$$\begin{aligned}
\int_{bU} \left| \int_U (Qf) \Phi_{x'}^\epsilon dt \right|^p dx' &\leq \int_{bU} \int_U |Qf|^p (\Phi_{x'}^\epsilon)^p dt dx' \\
&\leq \int_U |Qf|^p (1 - h_{\tau(\epsilon)}(-t_1)) \int_{bU} \psi_\epsilon(x' - t') dx' dt \\
&= \int_U |Qf|^p (1 - h_{\tau(\epsilon)}(-t_1)) dt
\end{aligned}$$

Here, $|Qf|^p (1 - h_{\tau(\epsilon)}(-t_1)) \leq |Qf|^p$ (which is in $L^1(U)$), and converges to 0 point-wise. Hence, the whole expression turns to 0 by Lebesgue's Theorem. For the second term, note that

$$|Q'_t \Phi_{x'}^\epsilon| = |(1 - h_{\tau(\epsilon)}(-t_1)) Q'_t \psi_\epsilon(x' - \cdot)| \lesssim \epsilon^{-n} (1 - h_{\tau(\epsilon)}(-t_1)).$$

Hence, we conclude:

$$\begin{aligned}
\int_{bU} \left| \int_U f(Q' \Phi_{x'}^\epsilon) dt \right|^p dx' &\leq \int_U |f|^p \left(\int_{bU} |Q' \Phi_{x'}^\epsilon|^p dx' \right) dt \\
&= \int_U |f|^p \left(\int_{B'_\epsilon(t')} |Q' \Phi_{x'}^\epsilon|^p dx' \right) dt \\
&\lesssim \frac{1}{\epsilon} \int_U |f|^p (1 - h_{\tau(\epsilon)}(-t_1)) dt.
\end{aligned}$$

Here now, for fixed $\epsilon > 0$, $\epsilon^{-1}|f|^p \in L^1(U)$,

$$\frac{1}{\epsilon} |f|^p (1 - h_{\tau(\epsilon)}(-t_1)) \leq \frac{1}{\epsilon} |f|^p,$$

and the left-hand side converges to 0 point-wise for $\tau \rightarrow 0$. So, by the Theorem of Lebesgue, there exists $\tau(\epsilon)$ such that

$$\frac{1}{\epsilon} \int_U |f|^p (1 - h_{\tau(\epsilon)}(-t_1)) dt \leq \epsilon.$$

This is our choice of $\tau(\epsilon)$ which has been left open before. So, we have just seen that

$$\lim_{\epsilon \rightarrow 0} (a_1(0, \cdot) f_b) * \psi_\epsilon(x') = \lim_{\epsilon \rightarrow 0} \int_U f(t) a_1(t) \phi_\epsilon(x - t) dt$$

in $L^p(bU)$. Recall that we had reduced the problem to showing (8). So, only

$$\lim_{\epsilon \rightarrow 0} \int_U f(t) (a_1(t) - a_1(x)) \phi_\epsilon(x - t) dt = 0$$

in $L^p(bU)$ remains to show. But, due to compactness, there exists a constant $C > 0$ such that $|a_1(t) - a_1(x)| \leq C\epsilon$ if $|t - x| \leq \epsilon$. Since $f \in L^p(U)$ and $|\phi_\epsilon| \leq 1$, the proof is finished easily. $\square$

We remark that the assumptions on the regularity of the boundary bM could be relaxed considerably.

4. BOUNDARY VALUES FOR THE $\bar{\partial}$ -OPERATOR

In this section, we will apply Friedrichs' extension lemma with boundary values, Theorem 3.3, to the $\bar{\partial}$ -operator. Recall the following definition of $\bar{\partial}$ -boundary values that is common in complex analysis:

Definition 4.1. *Let $D \subset\subset \mathbb{C}^n$ be a bounded domain with smooth boundary bD , and $f \in L^p_{0,q}(D)$ with $\bar{\partial}f \in L^p_{0,q+1}(D)$ in the sense of distributions for $1 \leq p < \infty$. Then, we say that f has weak $\bar{\partial}$ -boundary values $f_b \in L^p_q(bD)$ if*

$$\int_D \bar{\partial}f \wedge \phi + (-1)^q \int f \wedge \bar{\partial}\phi = \int_{bD} f_b \wedge \iota^*(\phi) \quad (9)$$

for all $\phi \in C^\infty_{n,n-q-1}(\bar{D})$, where $\iota : bD \hookrightarrow \mathbb{C}^n$ is the embedding of the boundary.

In fact, the left hand side of (9) depends only on the pull-back $\iota^*(\phi)$ of ϕ to bD , and so it defines a current on bD . Generally, this current is called the weak $\bar{\partial}$ -boundary value of f , and we say that f has got boundary values in L^p , if this current can be represented by an L^p -form as in Definition 4.1. See [H1] for a more detailed treatment of that topic. Boundary values as in Definition 4.1 are not uniquely defined. The reason is as follows: Let $r \in C^\infty(\mathbb{C}^n)$ be a defining function for D . So, $D = \{z \in \mathbb{C}^n : r(z) < 0\}$ and we may assume that $\|dr\| \equiv 1$ on bD . Then $\iota^*(dr) = 0$ implies $\iota^*(\bar{\partial}r) = -\iota^*(\partial r)$. Note that $\phi \in C^\infty_{n,n-q-1}(\bar{D})$ contains ∂r necessarily. Hence, $\iota^*(\bar{\partial}r) \wedge \iota^*(\phi) = 0$ on bD for all $\phi \in C^\infty_{n,n-q-1}(\bar{D})$. One should mention an example where weak $\bar{\partial}$ -boundary values occur:

Theorem 4.2. (Harvey-Polking [HP]) *Let r be the strictly plurisubharmonic defining function of a strictly pseudoconvex domain $D \subset\subset \mathbb{C}^n$, and $\omega \in L^1_{0,1}(D)$ with $\bar{\partial}\omega = 0$ and $|r|^{-1/2}\bar{\partial}r \wedge \omega \in L^1_{0,2}(D)$. Then there exists $f \in L^1(D)$ with weak $\bar{\partial}$ -boundary values $f_b \in L^1(bD)$ such that $\bar{\partial}f = \omega$.*

We will now show that Definition 4.1 is actually equivalent to Definition 2.2 if we make the right choices. So, let $M = \bar{D}$ with the underlying Riemannian structure on $\mathbb{C}^n$, $E = \Lambda^{0,q}T^*M$, $F = \Lambda^{0,q+1}T^*M$, and

$$Q := \bar{\partial} : C^\infty_{0,q}(M) = C^\infty(M, E) \longrightarrow C^\infty_{0,q+1}(M) = C^\infty(M, F).$$

Note that $Q^* = - * \partial *$. For $u, v \in C^\infty_{0,q}(\bar{D})$:

$$(u, v)_M = \int_M \langle u, v \rangle_E dV_{\mathbb{C}^n} = \int_M u \wedge * \bar{v}.$$

In order to reformulate (9), let $g := (-1)^{q+1} * \bar{\phi} \in C^\infty_{0,q+1}(D)$. Then $\phi = * \bar{g}$. So,

$$\int_D \bar{\partial}f \wedge \phi = \int_D \bar{\partial}f \wedge * \bar{g} = \int_M \langle \bar{\partial}f, g \rangle_F dV_M = (Qf, g)_M,$$

and

$$\begin{aligned} \int_D f \wedge \bar{\partial} \phi &= - \int_D f \wedge * * \bar{\partial} * * \phi = \int_D f \wedge * \overline{Q^* * \phi} \\ &= (-1)^{q+1} \int_M \langle f, Q^* g \rangle_E dV_{\mathbb{C}^n} = (-1)^{q+1} (f, Q^* g)_M. \end{aligned}$$

Hence, in the notation of Definition 2.2, the left hand side of (9) reads exactly as $(Qf, g)_M - (f, Q^*g)_M$. For the right hand side, recall that we have chosen the defining function r such that $\|dr\| \equiv 1$ on bD . That implies $dS_{bD} = \iota^*(dr)$. Note that there is a $(0, q)$ -form $f'_b \in C^\infty(bD, \Lambda^{0,q} T^* \mathbb{C}^n|_{bD})$ such that $\iota^*(f'_b) = f_b$. Since $\iota^* dr = 0$ and $dr \wedge \phi = \bar{\partial} r \wedge \phi$, we compute

$$\begin{aligned} f_b \wedge \iota^*(\phi) &= \iota^*(f'_b \wedge \phi) = \iota^*(([dr \wedge f'_b \wedge \phi]) * dr) = (*[\bar{\partial} r \wedge f'_b \wedge \phi]) dS_{bD} \\ &= (*[\bar{\partial} r \wedge f'_b \wedge * \bar{g}]) dS_{bD} = \langle \bar{\partial} r \wedge f'_b, g \rangle_F dS_{bM} \\ &= \langle Q(r f'_b), g \rangle_F dS_{bM} = \frac{1}{i} \langle \sigma_Q(\cdot, \nu^b) f'_b, g \rangle_F dS_{bM}. \end{aligned}$$

So, we have

$$\int_{bD} f_b \wedge \iota^*(\phi) = \frac{1}{i} \int_{bM} \langle \sigma_Q(x, \nu^b) f'_b, g \rangle_{F_x} dS_{bM}(x),$$

and recognize therefore:

Lemma 4.3. *$f \in L^p_{0,q}(D)$ with $\bar{\partial} f \in L^p_{0,q+1}(D)$ has weak $\bar{\partial}$ -boundary values $f_b \in L^p_q(bD)$ according to Definition 4.1 exactly if it has $\bar{\partial}$ -boundary values $f'_b \in L^p(bD, \Lambda^{0,q} T^* \mathbb{C}^n|_{bD})$ according to Definition 2.2.*

So, we are now in the position to translate Theorem 3.3 into the Friedrichs' extension lemma with boundary values for the $\bar{\partial}$ -operator:

Theorem 4.4. *Let $D \subset \subset \mathbb{C}^n$ be a bounded domain with smooth boundary bD , and $f \in L^p_{0,q}(D)$ with $\bar{\partial} f \in L^p_{0,q+1}(D)$ in the sense of distributions for $1 \leq p < \infty$. Then f has weak $\bar{\partial}$ -boundary values $f_b \in L^p_q(bD)$ according to Definition 4.1 exactly if there is a sequence $\{f_\epsilon\}$ in $C^\infty_{0,q}(\bar{D})$ such that $\lim_{\epsilon \rightarrow 0} f_\epsilon = f$ in $L^p_{0,q}(D)$, $\lim_{\epsilon \rightarrow 0} \bar{\partial} f_\epsilon = \bar{\partial} f$ in $L^p_{0,q+1}(D)$, and*

$$\lim_{\epsilon \rightarrow 0} \iota^*(f_\epsilon \wedge \phi) = f_b \wedge \iota^*(\phi) \quad \text{in } L^p_{2n-1}(bD)$$

for all $\phi \in C^\infty_{n,n-q-1}(\bar{D})$. If $r \in C^\infty(\mathbb{C}^n)$ is a defining function for D , i.e. $D = \{z \in \mathbb{C}^n : r(z) < 0\}$ and $dr \neq 0$ on bD , then the last condition is equivalent to

$$\lim_{\epsilon \rightarrow 0} \iota^*(f_\epsilon \wedge \partial r) = f_b \wedge \iota^*(\partial r) \quad \text{in } L^p_{q+1}(bD).$$

If $q = 0$, then this in turn is equivalent to

$$\lim_{\epsilon \rightarrow 0} \iota^*(f_\epsilon) = f_b \quad \text{in } L^p(bD).$$

5. REGULARITY OF THE BMK FORMULA

The characterization of weak $\bar{\partial}$ -boundary values by approximation is a quite useful tool because it allows us to simply work in the C^∞ -category in many situations. As an application, we will derive the Bochner-Martinelli-Koppelman formula for L^p -forms with weak $\bar{\partial}$ -boundary values. Before doing that, we present another technical but useful result. For convenience of the reader, let us recall shortly the Bochner-Martinelli-Koppelman formula.

Definition 5.1. *Let $0 \leq q \leq n$. The Bochner-Martinelli-Koppelman kernel B_{nq} in $\mathbb{C}^n$ is then given as*

$$B_{nq}(\zeta, z) = \frac{(n-1)!}{2^{q+1}\pi^n} \frac{1}{\|\zeta - z\|^{2n}} \sum_{\substack{j,J, \\ |L|=q+1}} \epsilon_{jJ}^L(\bar{\zeta}_j - \bar{z}_j)(*d\zeta^L) \wedge d\bar{z}^J,$$

where

$$\epsilon_B^A := \begin{cases} \text{sign } \pi & , \text{ if } A = B \text{ as sets and } \pi \text{ is a permutation with } B = \pi A, \\ 0 & , \text{ if } A \neq B. \end{cases}$$

Moreover, let $B_{n,-1} \equiv 0$.

Now, let $D \subset\subset \mathbb{C}^n$ be a bounded domain with C^1 -smooth boundary bD . If g is a measurable $(0, q+1)$ -form on D , let

$$\mathbf{B}_q^D g(z) := \int_D g(\zeta) \wedge B_{nq}(\zeta, z),$$

and if f is a measurable q -form on bD , let

$$\mathbf{B}_q^{bD} f(z) := \int_{bD} f(\zeta) \wedge B_{nq}(\zeta, z),$$

provided, the integrals do exist. Then:

Theorem 5.2. (BMK formula [K]) *Let $D \subset\subset \mathbb{C}^n$ be a bounded domain with C^1 -smooth boundary bD , $1 \leq q \leq n$, and $f \in C_{0,q}^1(\bar{D})$. Then:*

$$f(z) = \mathbf{B}_q^{bD} f(z) - \mathbf{B}_q^D(\bar{\partial}f)(z) - \bar{\partial}_z \mathbf{B}_{q-1}^D f(z), \quad (10)$$

where $\mathbf{B}_{q-1}^D f \in C_{0,q-1}^1(D)$.

In the following, we will show that (10) is still valid under the assumption that $f \in L_{0,q}^1(D)$ with $\bar{\partial}f \in L_{0,q+1}^1(D)$ has weak $\bar{\partial}$ -boundary values $f_b \in L_q^1(bD)$. It is well-known that

Lemma 5.3. *Let $D \subset\subset \mathbb{C}^n$ be a bounded domain. Then, $\mathbf{B}_q^D$ defines a bounded linear operator*

$$L_{0,q+1}^p(D) \rightarrow L_{0,q}^r(D)$$

for all $1 \leq p, r \leq \infty$ such that $1/r > 1/p - 1/(2n)$.

This is a direct consequence of $\|B_{nq}(\zeta, z)\| \lesssim \|\zeta - z\|^{2n-1}$ and Young's inequality, which is usually used for estimating integral operators (cf. for example [LM], Proposition III.5.35). In order to estimate the BMK boundary operator $\mathbf{B}_q^{bD}$, we need a more general version of such an inequality. So, we will make use of the following technical result. The proof can be found in [R1], Theorem 3.3.4:

Theorem 5.4. *Let $1 \leq t \leq s < \infty$ and $1 \leq a, b \leq \infty$ be fixed, (X, μ) and (Y, ν) measure spaces with $\mu(X) < \infty$ and $\nu(Y) < \infty$, and K a $\mu \times \nu$ -measurable function on $X \times Y$ such that*

$$\int_X |K(x, y)|^t d\mu(x) \leq g(y) \quad \text{for almost all } y \in Y, \quad (11)$$

$$\int_Y |K(x, y)|^s d\nu(y) \leq h(x) \quad \text{for almost all } x \in X, \quad (12)$$

where $g \in L^a(Y)$ and $h \in L^b(X)$. Then:

I. The linear operator $f \mapsto \mathbf{T}f$ which is given by

$$\mathbf{T}f(y) = \int_X K(x, y)f(x)d\mu(x)$$

for almost all $y \in Y$ defines a bounded operator $\mathbf{T} : L^p(X) \rightarrow L^r(Y)$ for all $1 \leq p, r \leq \infty$ satisfying

$$p \geq \begin{cases} \frac{t}{t-1} & , \text{ if } t > 1, \\ \infty & , \text{ if } t = 1, \end{cases} \quad (13)$$

and

$$r \leq at.$$

II. The mapping $f \mapsto \mathbf{T}f$ is bounded as an operator $\mathbf{T} : L^p(X) \rightarrow L^1(Y)$ for $1 \leq p < \infty$ with

$$p \geq \begin{cases} \frac{sb}{sb-1} & , \text{ if } 1 < sb < \infty, \\ 1 & , \text{ if } b = \infty. \end{cases} \quad (14)$$

III. If (14) is satisfied and $sb \neq t$, then $f \mapsto \mathbf{T}f$ defines a bounded operator $\mathbf{T} : L^p(X) \rightarrow L^r(Y)$ for all $1 \leq r \leq \infty$ with

$$\frac{1}{r} = \left(\frac{sb}{sb-t} \right) \left(\frac{1}{p} + \frac{1}{t} - 1 \right) \quad (15)$$

and

$$r \leq t \left(a \frac{s-t}{s} + 1 \right). \quad (16)$$

We have made the following conventions: In (15), let $1/r = 0$ if $r = \infty$. If $b = \infty$, then (15) has to be interpreted as $\frac{1}{r} = \frac{1}{p} + \frac{1}{t} - 1$. If $a = \infty$, then (16) reads as $r \leq \infty$.

It is now easy to deduce:

Lemma 5.5. *Let $D \subset\subset \mathbb{C}^n$ be a bounded domain with C^1 -smooth boundary bD . Then, $\mathbf{B}_q^{bD}$ is bounded as an operator*

$$L_q^p(bD) \rightarrow L_{0,q}^p(D)$$

for all $1 \leq p < \infty$.

Proof. We will apply Theorem 5.4 to the operator $\mathbf{B}_q^{bD}$. So, let $X = bD$, $Y = D$ and

$$|K(x, y)| = |B_{nq}(x, y)| \leq \frac{A}{|x - y|^{2n-1}},$$

where $A > 0$ is a constant that depends only on D , q and n . We choose $t = 1$. It is not hard to prove that there are constants $C_0(D) > 0$ and $C_1(D) > 0$ such that

$$\int_X |K(x, y)|^t d\mu(x) \leq C_0(D) + C_1(D) |\log \delta(y)| =: g(y),$$

where

$$\delta(y) := \text{dist}(y, bD).$$

For a proof, we refer to [R1], Lemma 3.3.1. It is easy to see that $|g|^a$ is integrable over $Y = D$ for all powers $1 \leq a < \infty$. So, we remark that $g \in L^a(Y)$ for all $1 \leq a < \infty$ (cf. [R1], Lemma 3.3.3). Now, choose $s > 1$ such that

$$1 = t < s < \frac{2n}{2n-1}.$$

Then

$$h(x) := \int_Y |K(x, y)|^s d\nu(y)$$

is uniformly bounded (independent of $x \in X$). Hence $h \in L^\infty(X)$. So, the assumptions of Theorem 5.4 are fulfilled for $X = bD$, $Y = D$, $\mathbf{T} = \mathbf{B}_q^{bD}$, $1 = t < s$, $h \in L^\infty(X)$, i.e. $b = \infty$, and $g \in L^a(Y)$ for all $1 \leq a < \infty$. We conclude that $\mathbf{B}_q^{bD}$ defines a bounded linear operator $\mathbf{B}_q^{bD} : L_q^p(bD) \rightarrow L_{0,q}^r(D)$ for all $1 \leq p, r < \infty$ such that

$$\frac{1}{r} = \frac{1}{p} + \frac{1}{t} - 1 = \frac{1}{p}.$$

□

We have now provided all the tools that are needed to derive the Bochner-Martinelli-Koppelman formula for L^p -forms with weak $\bar{\partial}$ -boundary values as an application of Friedrichs' extension theorem with boundary values. So, let $D \subset\subset \mathbb{C}^n$ be a bounded domain with smooth boundary, $1 \leq r, p < \infty$, $f \in L_{0,q}^p(D)$ with $\bar{\partial}f \in L_{0,q+1}^r(D)$ in the sense of distributions and assume that f has weak $\bar{\partial}$ -boundary values $f_b \in L_q^p(bD)$ according to Definition 4.1.

Then, by Theorem 4.4, it follows that there exists a sequence $\{f_\epsilon\}$ in $C_{0,q}^\infty(\overline{D})$ such that

$$\begin{aligned}\lim_{\epsilon \rightarrow 0} f_\epsilon &= f \text{ in } L_{0,q}^1(D), \\ \lim_{\epsilon \rightarrow 0} \bar{\partial} f_\epsilon &= \bar{\partial} f \text{ in } L_{0,q+1}^1(D),\end{aligned}$$

and

$$\lim_{\epsilon \rightarrow 0} f_\epsilon|_{bD} \wedge \iota^*(\phi) = \lim_{\epsilon \rightarrow 0} \iota^*(f_\epsilon \wedge \phi) = f_b \wedge \iota^*(\phi) \text{ in } L_{2n-1}^1(bD) \quad (17)$$

for all $\phi \in C_{n,n-q-1}^\infty(\overline{D})$, where $\iota : bD \rightarrow \mathbb{C}^n$ denotes the embedding. In the following, we will simply write f_ϵ instead of $f_\epsilon|_{bD}$.

Now, the classical BMK formula, Theorem 5.2, implies:

$$f_\epsilon(z) = \mathbf{B}_q^{bD} f_\epsilon(z) - \mathbf{B}_q^D(\bar{\partial} f_\epsilon)(z) - \bar{\partial}_z \mathbf{B}_{q-1}^D f_\epsilon(z)$$

for all $z \in D$, which we permute to:

$$\bar{\partial}_z \mathbf{B}_{q-1}^D f_\epsilon(z) = \mathbf{B}_q^{bD} f_\epsilon(z) - \mathbf{B}_q^D(\bar{\partial} f_\epsilon)(z) - f_\epsilon(z). \quad (18)$$

By Lemma 5.3 and Lemma 5.5, we know that the applications

$$\begin{aligned}\mathbf{B}_q^{bD} : L_q^1(bD) &\rightarrow L_{0,q}^1(D), \\ \mathbf{B}_q^D : L_{0,q+1}^1(D) &\rightarrow L_{0,q}^1(D)\end{aligned}$$

are continuous. Hence, the right hand side of (18) converges in $L_{0,q}^1(D)$ to a form

$$G = \mathbf{B}_q^{bD} f_b - \mathbf{B}_q^D(\bar{\partial} f) - f \in L_{0,q}^1(D). \quad (19)$$

To see this, note that the Bochner-Martinelli-Koppelman kernel $B_{nq}(\zeta, z)$ is a $(n, n-q-1)$ -form in ζ . So, (17) can be used.

Since

$$\begin{aligned}\lim_{\epsilon \rightarrow 0} \mathbf{B}_{q-1}^D f_\epsilon &= \mathbf{B}_{q-1}^D f \text{ in } L_{0,q-1}^1(D), \\ \lim_{\epsilon \rightarrow 0} \bar{\partial}_z \mathbf{B}_{q-1}^D f_\epsilon &= G \text{ in } L_{0,q}^1(D),\end{aligned}$$

it follows that G actually is the $\bar{\partial}$ -derivate in the sense of distributions:

$$G = \bar{\partial}_z \mathbf{B}_{q-1}^D f.$$

Applying Lemma 5.3 and Lemma 5.5 again, we observe that

$$\begin{aligned}\mathbf{B}_q^{bD} f_b &\in L_{0,q}^p(D), \\ \mathbf{B}_q^D(\bar{\partial} f) &\in L_{0,q}^r(D).\end{aligned}$$

So, the right hand side of (19), and therefore G , is in $L_{0,q}^r(D) \cap L_{0,q}^p(D)$.

We summarize:

Theorem 5.6. (BMK formula for L^p -forms) *Let $D \subset\subset \mathbb{C}^n$ be a bounded domain with smooth boundary and $0 \leq q \leq n$. Moreover, let $1 \leq r, p < \infty$ and $f \in L^p_{0,q}(D)$ with $\bar{\partial}f \in L^r_{0,q+1}(D)$, such that f has weak $\bar{\partial}$ -boundary values $f_b \in L^p_q(bD)$. Then*

$$\begin{aligned} \mathbf{B}_q^{bD} f_b &\in L^p_{0,q}(D), \\ \mathbf{B}_q^D(\bar{\partial}f) &\in L^r_{0,q}(D), \\ \mathbf{B}_{q-1}^D f &\in L^p_{0,q-1}(D) \cap \text{Dom}(\bar{\partial}), \\ \bar{\partial}\mathbf{B}_{q-1}^D f &\in L^r_{0,q}(D) \cap L^p_{0,q}(D), \end{aligned}$$

and

$$f(z) = \mathbf{B}_q^{bD} f_b(z) - \mathbf{B}_q^D(\bar{\partial}f)(z) - \bar{\partial}_z \mathbf{B}_{q-1}^D f(z)$$

for almost all $z \in D$.

REFERENCES

- [A] H. W. ALT, *Lineare Funktionalanalysis*, Springer-Verlag, Berlin, 1992.
- [BS] B. BERNDTSSON, N. SIBONY, The $\bar{\partial}$ -equation on a positive current, *Inventiones Math.* **147** (2002), 371–428.
- [F] K. FRIEDRICHS, The identity of weak and strong extensions of differential operators, *Trans. Amer. Math. Soc.* **55** (1944), 132–151.
- [HP] F. R. HARVEY, J. C. POLKING, The $\bar{\partial}$ -Neumann solution to the inhomogeneous Cauchy-Riemann equation in the ball in $\mathbb{C}^n$, *Trans. Amer. Math. Soc.* **281** (1984), 587–613.
- [H1] T. HEFER, Regularität von Randwerten der kanonischen Lösung der $\bar{\partial}$ -Gleichung auf streng pseudokonvexen Gebieten, *Bonner Math. Schr.* **320** (1999).
- [H2] T. HEFER, Boundary values of integral operators, *Math. Nachr.* **261/262** (2003), 85–104.
- [H3] L. HÖRMANDER, L^2 -estimates and existence theorems for the $\bar{\partial}$ -operator, *Acta Math.* **113** (1965), 89–152.
- [K] W. KOPPELMAN, The Cauchy integral for differential forms, *Bull. Amer. Math. Soc.* **73** (1967), 554–556.
- [LP] P. D. LAX, R. S. PHILLIPS, Local boundary conditions for dissipative symmetric linear differential operators, *Comm. Pure Appl. Math.* **13** (1960), 427–455.
- [LM] I. LIEB, J. MICHEL, *The Cauchy-Riemann Complex, Integral Formulae and Neumann Problem*, Vieweg, Braunschweig/Wiesbaden, 2002.
- [P] W. PARDON, The L^2 - $\bar{\partial}$ -cohomology of an algebraic surface, *Topology* **28** (1989), no. 2, 171–195.
- [PS1] W. PARDON, M. STERN, L^2 - $\bar{\partial}$ -cohomology of complex projective varieties, *J. Amer. Math. Soc.* **4** (1991), no. 3, 603–621.
- [PS2] W. PARDON, M. STERN, Pure Hodge structure on the L^2 -cohomology of varieties with isolated singularities, *J. reine angew. Math.* **533** (2001), 55–80.
- [R1] J. RUPPENTHAL, Zur Regularität der Cauchy-Riemannschen Differentialgleichungen auf komplexen Räumen, *Bonner Math. Schr.* **380** (2006).

- [R2] J. RUPPENTHAL, L^2 -theory for the $\bar{\partial}$ -operator on compact complex spaces, preprint 2009, *submitted*.
- [SCH] U. SCHULDENZUCKER, Regularität der kanonischen Lösung der Cauchy-Riemannschen Differentialgleichungen in Rand- L^p -Räumen, *Bonner Math. Schr.* **265** (1994).
- [S] H. SKODA, Valeurs au bord pour les solution de l'opérateur d'' , et caractérisation des zéros des fonctions de la classe de Nevanlinna, *Bull. Soc. Math. France* **104** (1976), 225–299.
- [T] M. E. TAYLOR, *Partial Differential Equations I*, Springer-Verlag, Berlin, 1996.

DEPARTMENT OF MATHEMATICS, UNIVERSITY OF WUPPERTAL, GAUSSSTR. 20, 42119 WUPPERTAL, GERMANY.

E-mail address: ruppenthal@uni-wuppertal.de