

A simplicial approach to the non-Abelian Chern-Simons path integral

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Motivation: Why is the theory of 3-manifold quantum invariants interesting?

- 1) *It is beautiful*: surprising relations between many different areas of mathematics/physics like
 - Algebra
 - low-dimensional Topology
 - Differential Geometry
 - Functional Analysis and Stochastic Analysis
 - Quantum field theory (in particular, Conformal field theory, Quantum Gravity, String theory)
- 2) *It is deep*: Fields Medals for Jones, Witten, Kontsevich
- 3) *It is useful*: Applications in Knot Theory and Quantum Gravity, ...

List of approaches to 3-manifold quantum invariants

Original heuristic approach

0. Chern-Simons path integrals approach (Witten)

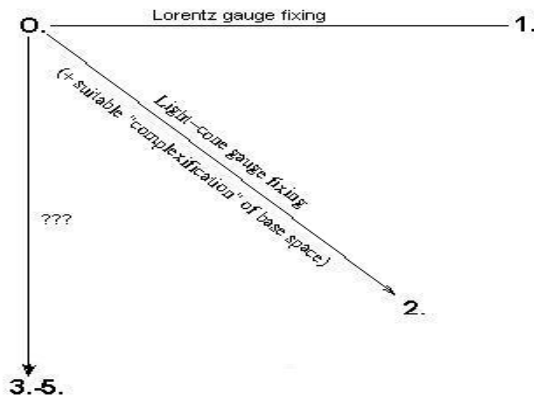
Rigorous perturbative approaches

1. Configuration space integrals
2. Kontsevich Integral

Rigorous non-perturbative approaches

- 3a. Quantum groups + Surgery (Reshetikhin/Turaev)
- 3b. Quantum groups + Shadow links (Turaev)
- 3c. Lattice gauge theories based on Quantum groups
4. Skein Modules
5. Geometric Quantization

Some relations between the approaches



Important open problems

- (P1) Chern-Simons path integral $\overset{??}{\longleftrightarrow}$ rigorous non-perturbative approaches 3a, 3b, 3c, 4, 5.
- (P2) Rigorous definition of original Chern-Simons path integral expressions?

Alternatively:

- (P2') Rigorous definition of Chern-Simons path integral expressions after suitable gauge fixing?

The shortterm goal

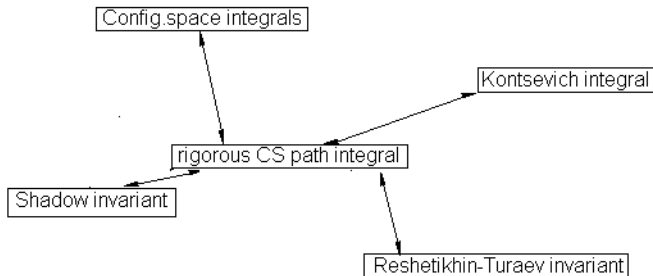
Our Aim

Make progress regarding (P1) and (P2')

Strategy

We consider special situation $M = \Sigma \times S^1$ and apply “torus gauge fixing” (cf. Blau/Thompson '93)

The longterm “goal” /dream



(all "arrows" being rigorous)

(or something analogous for BF_3 -theory with cosmological constant)

The Chern-Simons path integral

Fix

- M : oriented connected 3-manifold (usually compact)
- G : simply-connected simple Lie subgroup of $U(N)$ ($N \in \mathbb{N}$ fixed)
- $k \in \mathbb{R} \setminus \{0\}$ (usually $k \in \mathbb{N}$)

Space of gauge fields:

$$\mathcal{A} = \{A \mid A \text{ } \mathfrak{g}\text{-valued 1-form on } M\} \quad (\mathfrak{g} \subset \mathfrak{u}(N): \text{ Lie algebra of } G)$$

Action functional:

$$S_{CS} : \mathcal{A} \ni A \mapsto \frac{k}{4\pi} \int_M \text{Tr}(A \wedge dA + \frac{2}{3} A \wedge A \wedge A) \in \mathbb{R}$$

where $\text{Tr} := c \text{Tr}_{\text{Mat}(N, \mathbb{C})}$ for suitable normalisation constant $c \in \mathbb{R}$

Observation 1

S_{CS} is invariant under (orientation-preserving) diffeomorphisms

Fix

- “link” $L = (l_1, l_2, \dots, l_n)$, $n \in \mathbb{N}$, in M
- n -tuple $(\rho_1, \rho_2, \dots, \rho_n)$ of finite-dim. representations of G

“Definition”

$$Z(M, L) := \int \prod_i \text{Tr}_{\rho_i}(\text{Hol}_{l_i}(A)) \exp(iS_{CS}(A)) DA$$

where DA is the “Lebesgue measure” on \mathcal{A} and

$$\text{Hol}_{l_i}(A) := \lim_{n \rightarrow \infty} \prod_{k=1}^n \exp\left(\frac{1}{n} A(l'_i(\frac{k}{n}))\right) \quad (\text{“holonomy of } A \text{ around } l_i\text{”})$$

Observation 2

S_{CS} invariant under (orientation-preserving) diffeomorphisms $\Rightarrow Z(M, L)$ only depends on diffeomorphism class of M and isotopy class of L .

The shadow invariant for $M = \Sigma \times S^1$

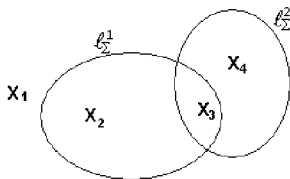
Special case: $M = \Sigma \times S^1$

Fix (framed) link $L = (l^1, l^2, \dots, l^n)$

Loop projections onto S^1 and Σ :

$$l_{S^1}^1, l_{S^1}^2, \dots, l_{S^1}^n \quad \text{and} \quad l_{\Sigma}^1, l_{\Sigma}^2, \dots, l_{\Sigma}^n$$

$D(L)$: graph in Σ generated by $l_{\Sigma}^1, l_{\Sigma}^2, \dots, l_{\Sigma}^n$



X_1, X_2, \dots, X_m : “faces” in $D(L)$

Gleams

Each X_t is equipped in a canonical way with a “gleam” $gl_t \in \mathbb{Z}$

Gleams $(gl_t)_t$ contain

- Information about crossings in $D(L)$,
- Information about winding numbers $\text{wind}(\mu_{S^1}^j)$

“Shadow of L ”

$sh(L) := (D(L), (gl_t)_t)$

Example 1: $D(L)$ has no crossing points

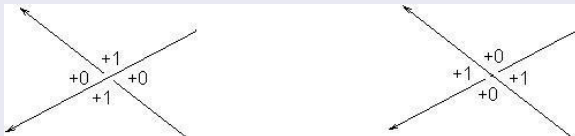
$$gl_t = \sum_{\{j \mid l_\Sigma^j \text{ touches } X_t\}} \text{wind}(l_{S^1}^j) \cdot \text{sgn}(X_t; l_\Sigma^j)$$

where

$$\text{sgn}(X_t; l_\Sigma^j) = \begin{cases} 1 & \text{if } X_t \text{ is "inside" of } l_\Sigma^j \\ -1 & \text{if } X_t \text{ is "outside" of } l_\Sigma^j \end{cases}$$

Example 2: $D(L)$ has crossing points but $\text{wind}(l_{S^1}^j) = 0, j \leq n$

Figure: Changes in the gleams at a given crossing point



Fix Cartan subalgebra \mathfrak{t} of \mathfrak{g} . Let k be as in Sec. 2 and additionally, $k > c_{\mathfrak{g}}$ ($c_{\mathfrak{g}}$ dual Coxeter number of \mathfrak{g}).

Colors and Colorings

- “color”: dominant weight of \mathfrak{g} (w.r.t. \mathfrak{t}) which is “integrable at level” $k - c_{\mathfrak{g}}$.
- \mathcal{C} : set of colors
- “link coloring” : mapping $\gamma : \{l_1, l_2, \dots, l_n\} \rightarrow \mathcal{C}$
- “area coloring”: mapping $\text{col} : \{X_1, \dots, X_m\} \rightarrow \mathcal{C}$.
- Col : set of area colorings

Fix link coloring $\gamma : \{l_1, l_2, \dots, l_n\} \rightarrow \mathcal{C}$.

Example 3: $\mathfrak{g} = \mathfrak{su}(2)$, \mathfrak{t} arbitrary

$$\mathcal{C} \cong \{0, 1, 2, \dots, k - 2\}$$

“Fusion coefficients” $N_{\alpha\beta}^{\gamma} \in \mathbb{N}_0$, $\alpha, \beta, \gamma \in \mathcal{C}$

$$N_{\alpha\beta}^{\gamma} = \sum_{\sigma \in W_k} (-1)^{\sigma} m_{\alpha}(\beta - \sigma(\gamma)),$$

where

- $m_{\alpha}(\beta)$: multiplicity of weight β in character of α .
- W_k : “quantum Weyl group” for \mathfrak{g} and k .

Remark

For our purposes the formula above is more useful than

$$N_{\alpha\beta}^{\gamma} = \sum_{\delta} \frac{S_{\alpha\delta} S_{\beta\delta} S_{\gamma^*\delta}}{S_{0\delta}} \quad \alpha, \beta, \gamma \in \mathcal{C},$$

where $(S_{\alpha\beta})_{\alpha\beta}$ is the S -matrix associated to \mathfrak{g} and k .

Some extra notation

- R_+ : set of positive real roots (w.r.t. fixed Weyl chamber)
- $\rho := \frac{1}{2} \sum_{\alpha \in R_+} \alpha$
- $\langle \cdot, \cdot \rangle$: Killing metric normalized such that $\langle \alpha, \alpha \rangle = 2$ if α is a long root.

“Shadow invariant” $|\cdot|$ for g and k

$$|L| = \sum_{\text{col} \in \text{Col}} \left(\prod_{i=1}^n N_{\gamma(l_i) \text{ col}(Y_i^+)}^{\text{col}(Y_i^-)} \right) \left(\prod_{t=1}^m (V_{\text{col}(X_t)})^{\chi(X_t)} (U_{\text{col}(X_t)})^{g|_t} \right) \\ \times \left(\prod_{p \in \text{CP}(L)} \text{symb}_q(\text{col}, p) \right) \quad \text{where}$$

$\chi(X_t)$: Euler characteristic of X_t

$Y_i^{+/-}$: face touching l_i^\pm from “inside” / “outside”

$$V_\lambda := \prod_{\alpha \in R_+} \frac{\sin \frac{\pi \langle \lambda + \rho, \alpha \rangle}{k}}{\sin \frac{\pi \langle \rho, \alpha \rangle}{k}}$$

$$U_\lambda := \exp\left(\frac{\pi i}{k} \langle \lambda, \lambda + 2\rho \rangle\right)$$

$\text{symb}_q(\text{col}, p)$: associated q-6j-symbol for $q := \exp(\frac{2\pi i}{k})$

Special case: $D(L)$ has no crossing points

$$|L| = \sum_{\text{col} \in \text{Col}} \left(\prod_{i=1}^n N_{\gamma(l_i) \text{ col}(Y_i^+)}^{\text{col}(Y_i^-)} \right) \left(\prod_{t=1}^m (V_{\text{col}(X_t)})^{\chi(X_t)} (U_{\text{col}(X_t)})^{\text{gl}_t} \right)$$

Aim of the rest of the Talk

Derive the formula above from the Chern-Simons path integral

From the CS path integral to the shadow invariant

Gauge group: $\mathcal{G} = C^\infty(M, G)$

\mathcal{G} operates on \mathcal{A} from the right by

$$A \cdot \Omega = \Omega^{-1} A \Omega + \Omega^{-1} d\Omega \quad \text{for } \Omega \in \mathcal{G}, A \in \mathcal{A}$$

Gauge Fixing: Choice of system \mathcal{A}_{gf} of representatives of \mathcal{A}/\mathcal{G}

Example: “Axial gauge fixing” for $M = \mathbb{R}^3$

In this case each $A \in \mathcal{A}$ can be written as $A = \sum_{i=0}^2 A_i dx_i$.

$$\mathcal{A}_{\text{gf}} = \{A \mid A_0 = 0\}$$

is “essentially” a gauge-fixing.

“Faddeev-Popov determinant”

If \mathcal{A}_{gf} is “nice” enough there is a function $\Delta_{\text{FadPop}} : \mathcal{A}_{\text{gf}} \rightarrow \mathbb{R}$ such that (informally)

$$\int_{\mathcal{A}} \chi(A) DA = \int_{\mathcal{A}_{\text{gf}}} \chi(A) \Delta_{\text{FadPop}}(A) DA|_{\mathcal{A}_{\text{gf}}}$$

for every \mathcal{G} -invariant function $\chi : \mathcal{A} \rightarrow \mathbb{C}$

Example

For $M = \mathbb{R}^3$ and $\mathcal{A}_{\text{gf}} := \mathcal{A}^\perp := \{A \in \mathcal{A} \mid A_0 = 0\}$ we have

$$\Delta_{\text{FadPop}}(A) = \text{const.} \quad \Rightarrow$$

$$\int_{\mathcal{A}} \chi(A) DA \sim \int_{\mathcal{A}^\perp} \chi(A^\perp) DA^\perp, \quad \text{with } DA^\perp := DA|_{\mathcal{A}^\perp}$$

Example for usefulness of applying a gauge fixing

For $M = \mathbb{R}^3$ and $\mathcal{A}_{\text{gf}} := \mathcal{A}^\perp := \{A \in \mathcal{A} \mid A_0 = 0\}$ we have

$$\begin{aligned} Z(M, L) &= \int_{\mathcal{A}} \prod_i \text{Tr}(\text{Hol}_{l_i}(A)) \exp(iS_{\text{CS}}(A)) DA \\ &\sim \int_{\mathcal{A}^\perp} \prod_i \text{Tr}(\text{Hol}_{l_i}(A)) \exp(iS_{\text{CS}}(A)) DA^\perp \\ &\stackrel{(*)}{=} \int_{\mathcal{A}^\perp} \prod_i \text{Tr}(\text{Hol}_{l_i}(A^\perp)) \exp(i\frac{k}{4\pi} \int \text{Tr}(dA^\perp \wedge A^\perp)) DA^\perp \end{aligned}$$

(*) holds because $A^\perp \wedge A^\perp \wedge A^\perp = 0$ for $A^\perp \in \mathcal{A}^\perp$.

The last integral involves a **Gauss-type measure!**

Torus Gauge

$$M = \mathbb{R}^3$$

$$A = \sum_{i=0}^2 A_i dx_i = A^\perp + A_0 dx_0 \text{ where } A^\perp := A_1 dx_1 + A_2 dx_2$$

$$\text{Note that } A^\perp \in \mathcal{A}^\perp := \{A \in \mathcal{A} \mid A(\frac{\partial}{\partial x_0}) = 0\}$$

$$\mathcal{A}_{\text{gf}} = \mathcal{A}^\perp = \{A \mid A_0 = 0\} \quad \textbf{Axial gauge fixing}$$

$$M = \Sigma \times S^1$$

$$\text{Each } A \in \mathcal{A} \text{ we have } A = A^\perp + A_0 dt \text{ with } A_0 \in C^\infty(\Sigma \times S^1, \mathfrak{g}) \text{ and } A^\perp \in \mathcal{A}^\perp := \{A \in \mathcal{A} \mid A(\frac{\partial}{\partial t}) = 0\}$$

$$dt \text{ and } \frac{\partial}{\partial t} \text{ obtained by lifting obvious 1-form/vector field on } S^1 \text{ to } \Sigma \times S^1$$

Problem: $\mathcal{A}_{\text{gf}} = \mathcal{A}^\perp$ is not a gauge fixing! We need larger space

1. Option

$$\mathcal{A}_{\text{gf}} = \mathcal{A}^\perp \oplus \{Bdt \mid B \in C^\infty(\Sigma, \mathfrak{g})\}$$

2. Option: “torus gauge” (Blau/Thompson '93)

$$\mathcal{A}_{\text{gf}} = \mathcal{A}^\perp \oplus \{Bdt \mid B \in C^\infty(\Sigma, \mathfrak{t})\}$$

where \mathfrak{t} is Lie algebra of fixed maximal torus $T \subset G$

Example: $T = \left\{ \begin{pmatrix} e^{i\theta} & 0 \\ 0 & e^{-i\theta} \end{pmatrix} \mid \theta \in \mathbb{R} \right\} \cong U(1)$ is a max. torus for $G = SU(2)$

Observation

$\Delta_{FadPop}(A^\perp + Bdt)$ only depends on $B \longrightarrow$ we set

$$\Delta_{FP}(B) := \Delta_{FadPop}(A^\perp + Bdt)$$

Example

For $G = SU(2)$ and T as above we have

$$\Delta_{FP}(B) \sim \prod_{\sigma} \sin^2(x(B(\sigma)))$$

for suitable isomorphism $x : \mathfrak{t} \rightarrow \mathbb{R}$

Not the full story!

Topological obstructions

- strictly speaking torus gauge is not a gauge
- we must allow gauge transformations which have a singularity in a fixed point $\sigma_0 \in \Sigma$ (this causes also A^\perp to have a singularity in σ_0)
- 1-1-correspondence

$$\{\text{relevant singularities of } A^\perp \text{ in } \sigma_0\} \longleftrightarrow [\Sigma, G/T] \cong \mathbb{Z}^{\dim(T)}$$

- extra summation $\sum_{h \in [\Sigma, G/T]} \cdots$ and factor depending on h in our formulas

Torus Gauge applied to CS theory on $M = \Sigma \times S^1$

“Definition” of $\Delta_{FP} \Rightarrow$

$$\begin{aligned}
 Z(M, L) &= \int_{\mathcal{A}} \prod_i \text{Tr}_{\rho_i}(\text{Hol}_{l_i}(A)) \exp(iS_{CS}(A)) DA \\
 &\sim \int_{C^\infty(\Sigma, \mathfrak{t})} \int_{\mathcal{A}^\perp} \prod_i \text{Tr}_{\rho_i}(\text{Hol}_{l_i}(A^\perp + Bdt)) \exp(iS_{CS}(A^\perp + Bdt)) DA^\perp \\
 &\quad \times \Delta_{FP}(B) DB \quad \text{where}
 \end{aligned}$$

- DA^\perp : “Lebesgue measure” on \mathcal{A}^\perp
- DB : “Lebesgue measure” on $C^\infty(\Sigma, \mathfrak{t})$

1. important Observation

$S_{CS}(A^\perp + Bdt)$ quadratic in A^\perp for fixed B

$\mathcal{A}_{\Sigma, V} :=$ Space of V -valued 1-forms on Σ for $V \in \{\mathfrak{g}, \mathfrak{t}, \mathfrak{t}^\perp\}$

Identification $\mathcal{A}^\perp \cong C^\infty(S^1, \mathcal{A}_{\Sigma, \mathfrak{g}})$

$$\mathcal{A}_c^\perp := \{A^\perp \in \mathcal{A}^\perp \mid A^\perp \text{ constant and } \mathcal{A}_{\Sigma, \mathfrak{t}\text{-valued}}\}$$

$$\check{\mathcal{A}}^\perp := \{A^\perp \in \mathcal{A}^\perp \mid \int_{S^1} A^\perp(t) dt \in \mathcal{A}_{\Sigma, \mathfrak{t}^\perp}\}$$

Decomposition $\mathcal{A}^\perp = \check{\mathcal{A}}^\perp \oplus \mathcal{A}_c^\perp$

2. important Observation

$$S_{CS}(\check{A}^\perp + A_c^\perp + Bdt) = S_{CS}(\check{A}^\perp + Bdt) + \frac{k}{2\pi} \int_{\Sigma} \text{Tr}(dA_c^\perp \cdot B)$$

Final heuristic integral formula

$$Z(M, L) \sim \int_{C^\infty(\Sigma, t)} \int_{\mathcal{A}_c^\perp} \int_{\check{A}^\perp} \prod_i \text{Tr}_{\rho_i}(\text{Hol}_{l_i}(\check{A}^\perp + A_c^\perp + Bdt)) d\check{\mu}_B^\perp(\check{A}^\perp) \\
\times \triangle_{FP}(B) Z(B) \exp(i \frac{k}{2\pi} \int_\Sigma \text{Tr}(dA_c^\perp \cdot B)) DA_c^\perp DB$$

where

$$d\check{\mu}_B^\perp(\check{A}^\perp) := \frac{1}{Z(B)} \exp(i S_{CS}(\check{A}^\perp + Bdt)) D\check{A}^\perp$$

with $Z(B) := \int \exp(i S_{CS}(\check{A}^\perp + Bdt)) D\check{A}^\perp$.

3. important Observation

Both $d\check{\mu}_B^\perp$ and $\exp(i \frac{k}{2\pi} \int_\Sigma \text{Tr}(dA_c^\perp \cdot B)) DA_c^\perp DB$ are of
 “Gauss-type”

Comment

Useful to fix an auxiliary Riemannian metric g_Σ on Σ :

→

- Hodge star operator $\star : \mathcal{A}_{\Sigma, g} \rightarrow \mathcal{A}_{\Sigma, g}$
- scalar product $\ll \cdot, \cdot \gg$ on $\mathcal{A}_{\Sigma, g}$

→

- operator $\star : C^\infty(S^1, \mathcal{A}_{\Sigma, g}) \rightarrow C^\infty(S^1, \mathcal{A}_{\Sigma, g})$
- scalar product $\ll \cdot, \cdot \gg$ on $C^\infty(S^1, \mathcal{A}_{\Sigma, g})$

Application: rewriting $d\check{\mu}_B^\perp = \exp(iS_{CS}(\check{A}^\perp + Bdt))D\check{A}^\perp$

$$d\check{\mu}_B^\perp = \exp(-i \frac{k}{4\pi} \ll \check{A}^\perp, \star(\frac{\partial}{\partial t} + \text{ad}(B))\check{A}^\perp \gg) D\check{A}^\perp$$

(Observe that $\star(\frac{\partial}{\partial t} + \text{ad}(B))$ is symmetric w.r.t. $\ll \cdot, \cdot \gg$)

Rigorous implementation

The “discretization approach”

Fix $m \in \mathbb{N}$ and fix triangulation of Σ , K being the underlying simplicial complex.

Let $C^p(K, V)$ denote the space of V -valued p -cochains for K , $p \in \mathbb{N}_0$

Discretization based on replacements

- $\mathcal{B} = C^\infty(\Sigma, \mathfrak{t}) = \Omega^0(\Sigma, \mathfrak{t}) \longrightarrow C^0(K, \mathfrak{t})$
- $\mathcal{A}_{\Sigma, \mathfrak{g}} = \Omega^1(\Sigma, \mathfrak{g}) \longrightarrow C^1(K, \mathfrak{g})$
- $\mathcal{A}^\perp \cong C^\infty(S^1, \mathcal{A}_{\Sigma, \mathfrak{g}}) \longrightarrow \text{Maps}(\mathbb{Z}_m, C^1(K, \mathfrak{g}))$

Remark

Apart from K we also use the dual K' of K in order to define a discrete Hodge star operator.

We need to use “field doubling” \rightarrow we are now considering BF_3 theory !
 (cf. Adams’ work on Simplicial Abelian Gauge Theories).

Evaluation of the path integral

Recall:

- We derived the heuristic formula

$$\begin{aligned} Z(M, L) &\sim \int_{C^\infty(\Sigma, \mathfrak{t})} \int_{\mathcal{A}_c^\perp} \int_{\check{A}^\perp} \prod_i \text{Tr}_{\rho_i}(\text{Hol}_{l_i}(\check{A}^\perp + A_c^\perp + Bdt)) d\check{\mu}_B^\perp(\check{A}^\perp) \\ &\quad \times \Delta_{FP}(B) Z(B) \exp(i \frac{k}{2\pi} \int_\Sigma \text{Tr}(dA_c^\perp \cdot B)) DA_c^\perp DB \end{aligned}$$

- We are able to make rigorous sense of the r.h.s. of $Z(M, L)$

Let us now restrict ourselves to the special case where L has no crossing points and evaluate $Z(M, L)$ explicitly.

For simplicity: heuristic treatment

Some properties of rigorous analogue of $\check{\mu}_B^\perp$

- oscillatory complex measure of “Gauss-type”
- normalized
- zero mean
- non-definite covariance operator

Toy model: complex “Gauss-type” measure μ on \mathbb{R}^2

$$\mu(x) = \frac{1}{2\pi} \exp(i\frac{1}{2}\langle x, Cx \rangle) dx \quad \text{where } C = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$$

Clearly, $\langle v, Cv \rangle = 0$ for $v = (1, 0)$

$$\Rightarrow \lim_{\epsilon \rightarrow 0} \int \langle x, v \rangle^n e^{-\epsilon |x|^2} d\mu(x) = 0 \quad \text{for all } n \in \mathbb{N}$$

$$\Rightarrow \lim_{\epsilon \rightarrow 0} \int \Phi(\langle x, v \rangle) e^{-\epsilon |x|^2} d\mu(x) = \Phi(0) \quad (\text{for all “sufficiently nice” entire analytic functions } \Phi : \mathbb{R} \rightarrow \mathbb{R})$$

1. Step: Perform $\int \cdots d\check{\mu}_B^\perp(\check{A}^\perp)$

$$\int_{\check{A}^\perp} \prod_j \text{Tr}_{\rho_j}(\text{Hol}_{l_j}(\check{A}^\perp + A_c^\perp + Bdt)) d\check{\mu}_B^\perp(\check{A}^\perp)$$

$$= \prod_j \text{Tr}_{\rho_j}(\text{Hol}_{l_j}(0 + A_c^\perp + Bdt)) = \prod_j \text{Tr}_{\rho_j}(\exp(\int_{l_j} A_c^\perp + \int_{l_j} Bdt))$$

→

$$\begin{aligned} Z(M, L) \sim & \int_{C^\infty(\Sigma, t)} \int_{\mathcal{A}_c^\perp} \prod_j \text{Tr}_{\rho_j}(\exp(\int_{l_j} A_c^\perp + \int_{l_j} Bdt)) \\ & \times \Delta_{FP}(B) Z(B) \exp(i \frac{k}{2\pi} \int_\Sigma \text{Tr}(dA_c^\perp \cdot B)) DA_c^\perp DB \end{aligned}$$

2. Step: Perform $\int \cdots DA_c^\perp$

Observe

$$\textcircled{1} \quad \text{Tr}_{\rho_j}(e^b) = \sum_{\alpha} m_{\rho_j}(\alpha) e^{i\alpha(b)} \quad \text{if } b \in \mathfrak{t}$$

$$\textcircled{2} \quad \int \text{Tr}(dA_c^\perp \cdot B) = \ll B, \star dA_c^\perp \gg$$

$$\textcircled{3} \quad \alpha\left(\int_{R_\Sigma^j} A_c^\perp\right) = \alpha\left(\int_{R_\Sigma^j} dA_c^\perp\right) = \ll \alpha \cdot 1_{R_\Sigma^j}, \star dA_c^\perp \gg$$

→

$$\begin{aligned} & \int \prod_j \text{Tr}_{\rho_j} \left(\exp\left(\int_{I_j} A_c^\perp + \int_{I_j} B dt\right) \exp\left(i \frac{k}{2\pi} \int_{\Sigma} \text{Tr}(dA_c^\perp \cdot B)\right) \right) DA_c^\perp \\ &= \sum_{\alpha_1, \dots, \alpha_n} \left(\prod_j m_{\rho_j}(\alpha_j) \right) [\dots] \int \exp\left(i \ll \frac{k}{2\pi} B - \sum_j \alpha_j 1_{R_\Sigma^j}, \star dA_c^\perp \gg\right) DA_c^\perp \\ &\sim \sum_{\alpha_1, \dots, \alpha_n} \left(\prod_j m_{\rho_j}(\alpha_j) \right) [\dots] \delta\left(B - \frac{2\pi}{k} \sum_j \alpha_j 1_{R_\Sigma^j}\right) \end{aligned}$$

3. Step: Perform $\int \cdots DB$

$$Z(M, L)$$

$$\sim \sum_{\alpha_1, \dots, \alpha_n} \int (\prod_j m_{\rho_j}(\alpha_j)) (\Delta_{FP}(B) Z(B)) (\exp(\dots)) \delta(B - \frac{2\pi}{k} \sum_j \alpha_j 1_{R_\Sigma^j}) DB$$

$$\stackrel{(*)}{=} \sum_{\{B = \frac{2\pi}{k} \sum_j 1_{R_\Sigma^j} \alpha_j\}} (\prod_j m_{\rho_j}(\alpha_j)) (\Delta_{FP}(B) Z(B)) (\exp(\dots))$$

$$= \dots$$

$$= \sum_{\text{col} \in \text{Col}} (\prod_{j=1}^n N_{\gamma(l_j) \text{col}(Y_j^+)}^{\text{col}(Y_j^-)}) (\prod_{t=1}^m (V_{\text{col}(X_t)})^{\chi(X_t)}) (\prod_{t=1}^m \exp(\mathfrak{gl}_t U_{\text{col}(X_t)}))$$

$$= |L|$$

(step $(*)$ is not quite the full story

Recall

Topological obstructions

- strictly speaking torus gauge is not a gauge
- we must allow A_c^\perp to have a singularity in fixed point σ_0 of Σ
- 1-1-correspondence

$$\{\text{relevant singularities of } A_c^\perp \text{ in } \sigma_0\} \longleftrightarrow [\Sigma, G/T] \cong \mathbb{Z}^{\dim(T)}$$

- extra summation $\sum_{h \in [\Sigma, G/T]} \cdots$ (plus a term depending on the “winding number” of h) in some formulas
- this extra summation (combined with a suitable application of the Poisson summation formula) does indeed lead to the correct expressions at the end of last slide

Open Questions

Question 1

What about the case where L does have crossing points:

Do we obtain quantum 6j-symbols?

Question 2

Discretization approach possible for original (= non-gauge fixed) path integral?

References

- ① V.G. Turaev. Shadow links and face models of statistical mechanics. *J. Diff. Geom.*, 36:35–74, 1992.
- ② M. Blau and G. Thompson. Derivation of the Verlinde Formula from Chern-Simons Theory and the G/G model. *Nucl. Phys.*, B408(1):345–390, 1993.
- ③ D. H. Adams. R-Torsion and Linking Numbers from Simplicial Abelian Gauge Theories. [arXiv:hep-th/9612009]
- ④ A. Hahn. An analytic Approach to Turaev's Shadow Invariant, *J. Knot Th. Ram.* 17(11):1327–1385, 2008 [arXiv:math-ph/0507040v7].
- ⑤ A. Hahn. From simplicial Chern-Simons theory to the shadow invariant I, 2012 [arXiv:1206.0439]
- ⑥ A. Hahn. From simplicial Chern-Simons theory to the shadow invariant II, 2012 [arXiv:1206.0441]